

This is a repository copy of *Cubic hypergeometric integrals of motion in affine Gaudin models*.

White Rose Research Online URL for this paper:

<https://eprints.whiterose.ac.uk/146456/>

Version: Accepted Version

Article:

Lacroix, Sylvain, Vicedo, Benoit orcid.org/0000-0003-3003-559X and Young, Charles (2020) Cubic hypergeometric integrals of motion in affine Gaudin models. *Advances in Theoretical and Mathematical Physics*. ISSN 1095-0761

<https://doi.org/10.4310/ATMP.2020.v24.n1.a5>

Reuse

Items deposited in White Rose Research Online are protected by copyright, with all rights reserved unless indicated otherwise. They may be downloaded and/or printed for private study, or other acts as permitted by national copyright laws. The publisher or other rights holders may allow further reproduction and re-use of the full text version. This is indicated by the licence information on the White Rose Research Online record for the item.

Takedown

If you consider content in White Rose Research Online to be in breach of UK law, please notify us by emailing eprints@whiterose.ac.uk including the URL of the record and the reason for the withdrawal request.

ADV. THEOR. MATH. PHYS.
Volume ?, Number ?, 1–32, ?

Cubic hypergeometric integrals of motion in affine Gaudin models

SYLVAIN LACROIX, BENOÎT VICEDO AND CHARLES YOUNG

We construct cubic Hamiltonians for quantum Gaudin models of affine types $\widehat{\mathfrak{sl}}_M$. They are given by hypergeometric integrals of a form we recently conjectured in [LVY]. We prove that they commute amongst themselves and with the quadratic Hamiltonians. We prove that their vacuum eigenvalues, and their eigenvalues for one Bethe root, are given by certain hypergeometric functions on a space of affine opers.

1. Introduction

The *quantum Gaudin model* [Gau14] can be defined for any symmetrizable Kac-Moody Lie algebra \mathfrak{g} . One chooses a collection z_1, \dots, z_N of distinct points in the complex plane and the model is defined by its quadratic Hamiltonians, which are the elements

$$\mathcal{H}_i = \sum_{\substack{j=1 \\ j \neq i}}^N \frac{\Xi^{(ij)}}{z_i - z_j}, \quad i = 1, \dots, N,$$

of the (suitably completed) tensor product $U(\mathfrak{g}^{\oplus N})$, where Ξ is the canonical element of $\mathfrak{g} \otimes \mathfrak{g}$ coming from the invariant bilinear form on \mathfrak{g} . In finite types \mathfrak{g} , these quadratic Hamiltonians are known to belong to a large commutative subalgebra $\mathcal{B} \subset U(\mathfrak{g}^{\oplus N})$ called the *Bethe* or *Gaudin* algebra, which is generated by the \mathcal{H}_i together with the central elements of $U(\mathfrak{g}^{\oplus N})$ and also (when $\text{rank}(\mathfrak{g}) > 1$) certain families of *higher Gaudin Hamiltonians* [FFR94, Fre05, Tal06, MTV06, Ryb08, Mol13, Ryb]. It is an interesting open question whether analogues of these higher Hamiltonians also exist when \mathfrak{g} is of affine type. Of particular interest is the case when \mathfrak{g} is of untwisted affine type since it is expected that quantum integrable field theories can be

described as affine Gaudin models associated with such Kac-Moody Lie algebras [FF11, Vic]. In a recent paper, [LVY], we conjectured that higher Hamiltonians do exist for \mathfrak{g} of untwisted affine type, and gave a broad conjecture for the form they should take as well as a precise conjecture for their eigenvalues on tensor products of irreducible highest-weight \mathfrak{g} -modules. According to [LVY] both the higher Hamiltonians and their eigenvalues should be given by integrals of hypergeometric type in the spectral plane (*i.e.* the copy of \mathbb{C} containing the marked points z_1, \dots, z_N).

In the present paper, we check the simplest case of those conjectures. The prediction is that there is a family of higher Hamiltonians for each *exponent* of \mathfrak{g} . Recall that, when \mathfrak{g} is of affine type, the exponents are a countably infinite (multi)set of integers whose pattern repeats modulo the Coxeter number (see *e.g.* [Kac90]). The quadratic Hamiltonians are associated with the exponent 1. The next case to check is that of Hamiltonians associated to cubic symmetric invariant tensors on the underlying finite-type Lie algebra. Such tensors in fact exist only in types A_{M-1} (*i.e.* \mathfrak{sl}_M) with $M \geq 3$. Equivalently, ${}^1A_{M-1}$ are the only untwisted affine types for which 2 is an exponent. In this paper we specialize to those types. We construct the cubic Hamiltonians and show that they commute amongst themselves and with the quadratic Hamiltonians. We also show that their eigenvalues are as predicted in [LVY] at least for Bethe vectors corresponding to 0 or 1 Bethe roots.

The paper is structured as follows.

In §2 we recall details of the affine algebra $\widehat{\mathfrak{sl}}_M$, its local completion and its vacuum Verma module. We also recall very briefly some concepts about vertex algebras.

In §3 we define the algebra of observables of the quantum Gaudin model, and define the states $\varsigma_1(z)$ and $\varsigma_2(z)$ used to construct the quadratic and cubic Hamiltonians respectively. The main technical results of the paper are then Theorems 3.3 and 3.4. The former shows that the $(\cdot)_0$ th products (in the sense of vertex algebras) between these states vanish modulo certain *twisted derivatives* and *translates*. The latter shows that the same is true of the action of the diagonal copy of $\widehat{\mathfrak{sl}}_M$ on $\varsigma_i(z)$, $i = 1, 2$.

These statements allow us to prove our main result, Theorem 3.8, which establishes that the quadratic and cubic Hamiltonians commute amongst themselves and with the diagonal action of $\widehat{\mathfrak{sl}}_M$. The Hamiltonians \hat{Q}_i^γ , $i = 1, 2$, are defined in (3.14): the superscript γ denotes a Pochhammer contour in $\mathbb{C} \setminus \{z_1, \dots, z_N\}$.

Cubic hypergeometric integrals of motion in affine Gaudin models 3

(There is a slight subtlety because the \hat{Q}_1^γ are not exactly the standard quadratic Hamiltonians \mathcal{H}_j . Thus, we also show in Theorem 3.8 that the \hat{Q}_i^γ , $i = 1, 2$, commute with the \mathcal{H}_j . We do so using another result, Theorem 3.5.)

In §4 we recall the conjectured eigenvalues for the Hamiltonians \hat{Q}_i^γ , $i = 1, 2$, from [LVY]. Namely, the eigenvalues are obtained by putting a certain *affine oper* (coming from an affine *Miura oper*) into *quasi-canonical form*, and then integrating the resulting coefficient functions $v_i(z)$, $i = 1, 2$, along the same contour γ .

In §5 we check that the predicted eigenvalues of the cubic Hamiltonians are correct (for the quadratic Hamiltonians see [LVY]) for Bethe vectors with 0 and 1 Bethe roots, *i.e.* for the vacuum state and for Bethe eigenstates at one step down in the principal gradation.

Finally, in §6 we consider the special case of only $N = 2$ marked points. In that case we show that our Hamiltonians \hat{Q}_i^γ , $i = 1, 2$ coincide up to rescaling with the zero modes of, respectively, the GKO coset conformal vector ω , [GKO86], and a state W constructed in [BBSS88]. It is known that, after certain further specializations, these generate a copy of the W_3 algebra. This provides an interesting arena for checking conjectures about higher Hamiltonians, since higher integrals of motion of the W_3 algebra are already known.

Acknowledgements. CY is grateful to E. Mukhin for interesting discussions. SL thanks F. Delduc and M. Magro for interesting discussions. This work is partially supported by the French Agence Nationale de la Recherche (ANR) under grant ANR-15-CE31-0006 DefIS.

2. Vacuum verma modules for $\hat{\mathfrak{sl}}_M$

2.1. Loop realization of $\hat{\mathfrak{sl}}_M$

We work over \mathbb{C} . Pick and fix an integer $M \geq 3$. Let $\mathfrak{sl}_M[t, t^{-1}] = \mathfrak{sl}_M \otimes \mathbb{C}[t, t^{-1}]$ denote the Lie algebra of Laurent polynomials, in a formal variable t , with coefficients in the finite-dimensional simple Lie algebra \mathfrak{sl}_M . The Lie bracket on $\mathfrak{sl}_M[t, t^{-1}]$ is given by $[a \otimes f(t), b \otimes g(t)] := [a, b] \otimes f(t)g(t)$ for any $a, b \in \mathfrak{sl}_M$ and $f(t), g(t) \in \mathbb{C}[t, t^{-1}]$. Let $(\cdot | \cdot) : \mathfrak{sl}_M \times \mathfrak{sl}_M \rightarrow \mathbb{C}$ be the standard bilinear form on \mathfrak{sl}_M . It is given by

$$(2.1) \quad (X|Y) := \text{tr}(XY)$$

4 Sylvain Lacroix, Benoît Vicedo and Charles Young

where tr denotes the trace in the defining $M \times M$ matrix representation. The affine algebra $\widehat{\mathfrak{sl}}_M$ is the central extension of $\mathfrak{sl}_M[t, t^{-1}]$ by a one dimensional centre $\mathbb{C}\mathbf{k}$,

$$0 \longrightarrow \mathbb{C}\mathbf{k} \longrightarrow \widehat{\mathfrak{sl}}_M \longrightarrow \mathfrak{sl}_M[t, t^{-1}] \longrightarrow 0,$$

whose commutation relations are given by $[\mathbf{k}, \cdot] = 0$ and

$$[a \otimes f(t), b \otimes g(t)] := [a, b] \otimes f(t)g(t) - (\text{res}_t f dg)(a|b)\mathbf{k}.$$

Let $a_n := a \otimes t^n$ for $a \in \mathfrak{sl}_M$ and $n \in \mathbb{Z}$. The commutation relations can equivalently be written as

$$[a_m, b_n] = [a, b]_{n+m} - n\delta_{n+m,0}(a|b)\mathbf{k}.$$

Define the Lie algebra

$$\mathfrak{g} := \widehat{\mathfrak{sl}}_M \oplus \mathbb{C}\mathbf{d},$$

by declaring that the *derivation element* \mathbf{d} obeys $[\mathbf{d}, \mathbf{k}] = 0$ and $[\mathbf{d}, a \otimes f(t)] = a \otimes t\partial_t f(t)$ for all $a \in \mathfrak{sl}_M$ and $f(t) \in \mathbb{C}[t, t^{-1}]$.

2.2. Kac-Moody data

Recall, for example from [Kac90], that the Lie algebra \mathfrak{g} is isomorphic to the Kac-Moody algebra over \mathbb{C} of type ${}^1\mathbf{A}_{M-1}$. The Cartan matrix is $A := (a_{ij})_{i,j=0}^{M-1} = (2\delta_{ij} - \delta_{i+1,j} - \delta_{i-1,j})_{i,j=0}^{M-1}$, where addition of indices is modulo M . Fix a Cartan decomposition $\mathfrak{g} = \mathfrak{n}_- \oplus \mathfrak{h} \oplus \mathfrak{n}_+$ and sets of Chevalley-Serre generators $\{e_i\}_{i=0}^{M-1} \subset \mathfrak{n}_+$, $\{f_i\}_{i=0}^{M-1} \subset \mathfrak{n}_-$. Let $\{\check{\alpha}_i\}_{i=0}^{M-1} \subset \mathfrak{h}$ be the simple coroots of \mathfrak{g} and $\{\alpha_i\}_{i=0}^{M-1} \subset \mathfrak{h}^*$ the simple roots. They obey $a_{ij} = \langle \alpha_j, \check{\alpha}_i \rangle$ for $i, j \in \{0, 1, \dots, M-1\}$, where $\langle \cdot, \cdot \rangle : \mathfrak{h}^* \times \mathfrak{h} \rightarrow \mathbb{C}$ is the canonical pairing of the Cartan subalgebra \mathfrak{h} and its dual \mathfrak{h}^* . The defining relations of \mathfrak{g} are then

$$(2.2a) \quad [x, e_i] = \langle \alpha_i, x \rangle e_i, \quad [x, f_i] = -\langle \alpha_i, x \rangle f_i,$$

$$(2.2b) \quad [x, x'] = 0, \quad [e_i, f_j] = \check{\alpha}_i \delta_{i,j},$$

for any $x, x' \in \mathfrak{h}$ and $i, j \in \{0, 1, \dots, M-1\}$, together with the Serre relations

$$(2.2c) \quad (\text{ad } e_i)^{1-a_{ij}} e_j = 0, \quad (\text{ad } f_i)^{1-a_{ij}} f_j = 0.$$

The centre of \mathfrak{g} is one dimensional and spanned by the *central element* $\mathbf{k} := \sum_{i=0}^{M-1} \check{\alpha}_i$. The Cartan subalgebra has a basis consisting of the simple

Cubic hypergeometric integrals of motion in affine Gaudin models 5

coroots $\{\check{\alpha}_i\}_{i=0}^{M-1}$ together with the derivation element \mathbf{d} , which obeys

$$\langle \alpha_i, \mathbf{d} \rangle = \delta_{i,0}.$$

(This condition fixes \mathbf{d} uniquely up to the addition of a multiple of \mathbf{k} .)

The matrix $(a_{ij})_{i,j=1}^{M-1}$ obtained by removing the zeroth row and column of A is the Cartan matrix of finite type A_{M-1} , with \mathfrak{sl}_M the corresponding finite-dimensional simple Lie algebra. We can identify \mathfrak{sl}_M with the subalgebra of \mathfrak{g} generated by $\{e_i\}_{i=1}^{M-1} \subset \mathfrak{n}_+$ and $\{f_i\}_{i=1}^{M-1} \subset \mathfrak{n}_-$, and then the defining relations are as above. Let $\dot{\mathfrak{h}} := \text{span}_{\mathbb{C}}\{\check{\alpha}_i\}_{i=1}^{M-1} \subset \mathfrak{sl}_M$ denote its Cartan subalgebra and $\dot{\mathfrak{h}}^* = \text{span}_{\mathbb{C}}\{\alpha_i\}_{i=1}^{M-1}$ its dual. Let $\{\omega_i\}_{i=1}^{M-1} \subset \dot{\mathfrak{h}}^*$ and $\{\check{\omega}_i\}_{i=1}^{M-1} \subset \dot{\mathfrak{h}}$ be respectively the fundamental weights and coweights of \mathfrak{sl}_M .

2.3. Local completion and vacuum Verma module

For any $k \in \mathbb{C}$, let $U_k(\widehat{\mathfrak{sl}}_M)$ denote quotient of the enveloping algebra of $\widehat{\mathfrak{sl}}_M$ by the two-sided ideal generated by $\mathbf{k} - k$. For each $n \in \mathbb{Z}_{\geq 0}$ define the left ideal $J_n := U_k(\widehat{\mathfrak{sl}}_M) \cdot (\mathfrak{sl}_M \otimes t^n \mathbb{C}[t])$. That is, J_n is the linear span of monomials of the form $a_p \dots b_q c_r$ with $r \geq n$, $c \in \mathfrak{sl}_M$, and some non-negative number of elements $a, \dots, b \in \mathfrak{sl}_M$ and mode numbers $p, \dots, q \in \mathbb{Z}$. The inverse limit $\tilde{U}_k(\widehat{\mathfrak{sl}}_M) := \varprojlim U_k(\widehat{\mathfrak{sl}}_M)/J_n$ is a complete topological algebra, called the *local completion of $U(\widehat{\mathfrak{sl}}_M)$ at level k* . By definition, elements of $\tilde{U}_k(\widehat{\mathfrak{sl}}_M)$ are (possibly infinite) sums $\sum_{m \geq 0} X_m$ of elements $X_m \in U_k(\widehat{\mathfrak{sl}}_M)$ which truncate to finite sums when one works modulo any J_n , i.e. for every n , $X_m \in J_n$ for all sufficiently large m .

A module \mathcal{M} over $\widehat{\mathfrak{sl}}_M$ is called *smooth* if, for all $a \in \mathfrak{sl}_M$ and all $v \in \mathcal{M}$, $a_n v = 0$ for all sufficiently large n . A module \mathcal{M} has *level k* if $\mathbf{k} - k$ acts as zero on \mathcal{M} . Any smooth module of level k over $\widehat{\mathfrak{sl}}_M$ is also a module over $\tilde{U}_k(\widehat{\mathfrak{sl}}_M)$.

We have the Lie subalgebra $\mathfrak{sl}_M[t] \oplus \mathbb{C}\mathbf{k} \subset \widehat{\mathfrak{sl}}_M$. Let $\mathbb{C}|0\rangle^k$ denote the one-dimensional representation of this Lie algebra given by $(\mathbf{k} - k)|0\rangle^k = 0$ and $a_n|0\rangle^k = 0$ for all $n \geq 0$ and all $a \in \mathfrak{sl}_M$. Let \mathbb{V}_0^k denote the induced $\widehat{\mathfrak{sl}}_M$ -module:

$$\mathbb{V}_0^k = U(\widehat{\mathfrak{sl}}_M) \otimes_{U(\mathfrak{sl}_M[t] \oplus \mathbb{C}\mathbf{k})} \mathbb{C}|0\rangle^k.$$

This module \mathbb{V}_0^k is called the *vacuum Verma module of level k* . It is a smooth module. Concretely, \mathbb{V}_0^k is the linear span of vectors of the form $a_p \dots b_q |0\rangle^k$ with $a, \dots, b \in \mathfrak{sl}_M$ and strictly negative mode numbers $p, \dots, q \in \mathbb{Z}_{\leq -1}$. Elements of \mathbb{V}_0^k are called *states*.

6 Sylvain Lacroix, Benoît Vicedo and Charles Young

Let $[T, \cdot]$ be the derivation on $U_k(\widehat{\mathfrak{sl}}_M)$ defined by $[T, a_n] := -na_{n-1}$ and $[T, 1] := 0$. By setting $T(X|0)^k := [T, X]|0^k$ for any $X \in U_k(\widehat{\mathfrak{sl}}_M)$, one can then regard T also as a linear map $\mathbb{V}_0^k \rightarrow \mathbb{V}_0^k$, called the *translation operator*.

2.4. Vertex algebra structure

For every state $A \in \mathbb{V}_0^k$ and every $n \in \mathbb{Z}$, there is an element $A_{(n)} \in \tilde{U}_k(\widehat{\mathfrak{sl}}_M)$, the n^{th} *formal mode* of A . These modes can be arranged into a formal power series

$$(2.3) \quad Y[A, u] := \sum_{n \in \mathbb{Z}} A_{(n)} u^{-n-1}$$

where $Y[\cdot, u]$ is called the *formal state-field map*. Their definition is as follows.

First, if $A = a_{-1}|0^k$ for some $a \in \mathfrak{sl}_M$ then $A_{(n)} := a_n$ for all $n \in \mathbb{Z}$, *i.e.*

$$Y[a_{-1}|0^k, u] := \sum_{n \in \mathbb{Z}} a_n u^{-n-1}.$$

Next, for all states $A, B \in \mathbb{V}_0^k$,

$$Y[TA, u] := \partial_u Y[A, u] \quad \text{and} \quad Y[A_{(-1)}B, u] := :Y[A, u]Y[B, u]:$$

where

$$(2.4) \quad :Y[A, u]Y[B, u]: := \left(\sum_{m < 0} A_{(m)} u^{-m-1} \right) Y[B, u] + Y[B, u] \left(\sum_{m \geq 0} A_{(m)} u^{-m-1} \right)$$

is called the *normal ordered product*. These assignments together recursively define $Y[C, u]$ for all $C \in \mathbb{V}_0^k$ (by writing $C = a_{-n}B$ for some $a \in \mathfrak{sl}_M$, $n \in \mathbb{Z}_{\geq 1}$ and $B \in \mathbb{V}_0^k$).

The *state-field map* $Y(\cdot, u) : \mathbb{V}_0^k \rightarrow \text{Hom}(\mathbb{V}_0^k, \mathbb{V}_0^k((u)))$ is obtained from $Y[A, u]$ by sending each formal mode to its image in $\text{End}(\mathbb{V}_0^k)$. This map $Y(\cdot, u)$ obeys a collection of axioms that make \mathbb{V}_0^k into a *vertex algebra*; see *e.g.* [FBZ04].

3. Quadratic and cubic Hamiltonians

3.1. The algebra of observables

Let $\mathbf{k} := (k_i)_{i=1}^N$ be a collection of $N \in \mathbb{Z}_{\geq 1}$ complex numbers $k_i \neq -M$ for $i = 1, \dots, N$. Consider the tensor product

$$\mathbb{V}_0^{\mathbf{k}} := \mathbb{V}_0^{k_1} \otimes \dots \otimes \mathbb{V}_0^{k_N}$$

of vacuum Verma modules. We can regard it as a module over the direct sum $\widehat{\mathfrak{sl}}_M^{\oplus N}$ of N copies of $\widehat{\mathfrak{sl}}_M$. Let $A^{(i)} \in \widehat{\mathfrak{sl}}_M^{\oplus N}$ denote the copy of $A \in \widehat{\mathfrak{sl}}_M$ in the i^{th} direct summand. Let $\mathbb{C}|0\rangle^{\mathbf{k}}$ denote the one-dimensional representation of the Lie subalgebra $(\mathfrak{sl}_M[t] \oplus \mathbb{C}\mathbf{k})^{\oplus N} \subset \widehat{\mathfrak{sl}}_M^{\oplus N}$ defined by $(\mathbf{k}^{(i)} - k_i)|0\rangle^{\mathbf{k}} = 0$ and $a_n^{(i)}|0\rangle^{\mathbf{k}} = 0$, for all $n \geq 0$, all $a \in \mathfrak{sl}_M$, and all $i \in \{1, \dots, N\}$. Then $\mathbb{V}_0^{\mathbf{k}}$ is the induced $\widehat{\mathfrak{sl}}_M^{\oplus N}$ -module:

$$\mathbb{V}_0^{\mathbf{k}} = U(\widehat{\mathfrak{sl}}_M^{\oplus N}) \otimes_{U((\mathfrak{sl}_M[t] \oplus \mathbb{C}\mathbf{k})^{\oplus N})} \mathbb{C}|0\rangle^{\mathbf{k}}.$$

Let $U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ denote the quotient of $U(\widehat{\mathfrak{sl}}_M^{\oplus N})$ by the two-sided ideal generated by $\mathbf{k}^{(i)} - k_i$ for all $i \in \{1, \dots, N\}$. We have the isomorphism

$$U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N}) \cong U_{k_1}(\widehat{\mathfrak{sl}}_M) \otimes \dots \otimes U_{k_N}(\widehat{\mathfrak{sl}}_M),$$

which sends $A^{(i)} \in \widehat{\mathfrak{sl}}_M^{\oplus N} \subset U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ to $1^{\otimes i-1} \otimes A \otimes 1^{\otimes N-i}$. Let J_n^N denote the left ideal in $U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ generated by $a_r^{(i)}$ for all $r \geq n$, $a \in \mathfrak{sl}_M$ and all $i \in \{1, \dots, N\}$. Let $\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N}) := \varprojlim U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})/J_n^N$ denote the inverse limit. It is a complete topological algebra, and

$$\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N}) \cong \tilde{U}_{k_1}(\widehat{\mathfrak{sl}}_M) \hat{\otimes} \dots \hat{\otimes} \tilde{U}_{k_N}(\widehat{\mathfrak{sl}}_M),$$

where $\hat{\otimes}$ denotes the completed tensor product. We call $\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ the *algebra of observables* of the Gaudin model.

We have the formal state-field map $Y[\cdot, u] : \mathbb{V}_0^{\mathbf{k}} \rightarrow \tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})[[u, u^{-1}]]$ and translation operator $T \in \text{End}(\mathbb{V}_0^{\mathbf{k}})$ defined as above. $\mathbb{V}_0^{\mathbf{k}}$ is a module over $\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$, and we get the state-field map $Y(\cdot, u) : \mathbb{V}_0^{\mathbf{k}} \rightarrow \text{Hom}(\mathbb{V}_0^{\mathbf{k}}, \mathbb{V}_0^{\mathbf{k}}((u)))$.

Let $\Delta : \widehat{\mathfrak{sl}}_M \hookrightarrow \widehat{\mathfrak{sl}}_M^{\oplus N}$ denote the diagonal embedding of $\widehat{\mathfrak{sl}}_M$ into $\widehat{\mathfrak{sl}}_M^{\oplus N}$,

$$\Delta x = \sum_{i=1}^N x^{(i)}$$

8 Sylvain Lacroix, Benoît Vicedo and Charles Young

It extends uniquely to an embedding $\Delta : U(\widehat{\mathfrak{sl}}_M) \hookrightarrow U(\widehat{\mathfrak{sl}}_M^{\oplus N}) = U(\widehat{\mathfrak{sl}}_M)^{\otimes N}$ (the usual N -fold coproduct). For all $a, b \in \mathfrak{sl}_M$ and $m, n \in \mathbb{Z}$,

$$[\Delta a_m, \Delta b_n] = \Delta[a, b]_{n+m} - n\delta_{n+m,0}(a|b) \sum_{i=1}^n \mathbf{k}^{(i)}$$

in $U(\widehat{\mathfrak{sl}}_M^{\oplus N})$. Therefore Δ descends to an embedding of the quotients, $\Delta : U_{|\mathbf{k}|}(\widehat{\mathfrak{sl}}_M) \hookrightarrow U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$, where $|\mathbf{k}| = \sum_{i=1}^N k_i$, and thence of their completions,

$$\Delta : \tilde{U}_{|\mathbf{k}|}(\widehat{\mathfrak{sl}}_M) \hookrightarrow \tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N}).$$

3.2. Invariant tensors on \mathfrak{sl}_M

Let I^a , $a = 1, \dots, \dim(\mathfrak{sl}_M)$, be a basis of \mathfrak{sl}_M and I_a , $a = 1, \dots, \dim(\mathfrak{sl}_M)$ the dual basis with respect to the non-degenerate bilinear form $(\cdot|\cdot)$ from (2.1). Let f^{ab}_c denote the structure constants,

$$(3.1) \quad [I^a, I^b] = f^{ab}_c I^c := \sum_c f^{ab}_c I^c,$$

where, from now on, we employ summation convention over repeated pairs of Lie algebra indices. Since the symmetric bilinear form $(\cdot|\cdot)$ is non-degenerate, we may and shall suppose that the basis is chosen such that

$$(I_a|I_b) = \delta_{ab}.$$

Then $I_a = I^a$ for each a and we no longer need to distinguish between upper and lower indices, and we shall write for example $f_{abc} = f^{ab}_c$. We have

$$\delta_{ab} = \text{tr } I_a I_b, \quad f_{abc} = \text{tr}[I_a, I_b] I_c = \text{tr}(I_a I_b I_c - I_b I_a I_c).$$

Recall that in general a tensor $t : \mathfrak{sl}_M \times \dots \times \mathfrak{sl}_M \rightarrow \mathbb{C}$ is *invariant* if

$$t([a, x], y, \dots, z) + t(x, [a, y], \dots, z) + \dots + t(x, y, \dots, [a, z]) = 0$$

for all x, y, \dots, z and a in \mathfrak{sl}_M . In particular, δ_{ab} and f_{abc} are the components of respectively a symmetric second-rank invariant tensor and a totally skew-symmetric third-rank invariant tensor.

The reason for specialising to \mathfrak{sl}_M in this paper is that in these types, and in no others, there exists a nonzero totally symmetric third-rank tensor.

Cubic hypergeometric integrals of motion in affine Gaudin models 9

It is unique up to normalization, and is given by $t(x, y, z) := \text{tr}(xyz + yxz)$. Let t_{abc} denote its components,

$$t_{abc} := t(I_a, I_b, I_c).$$

The Coxeter number and dual Coxeter number of \mathfrak{sl}_M (and of $\widehat{\mathfrak{sl}}_M$) are equal to M .

Lemma 3.1. *We have the following tensor identities*

$$\begin{aligned} f_{ade}f_{bdg}t_{ceg} &= -Mt_{abc} & t_{abc}t_{dbc} &= 2\frac{M^2-4}{M}\delta_{ad}, \\ f_{abc}f_{dbc} &= -2M\delta_{ad}, & f_{ade}t_{bgd}t_{ceg} &= \frac{M^2-4}{M}f_{abc}, \end{aligned}$$

and the tensors $f_{ade}f_{dfg}f_{ehi}t_{bfh}t_{cgi}$, $f_{acd}f_{def}f_{egh}f_{fij}t_{bgi}t_{chj}$ and $f_{acd}f_{def}f_{egh}f_{fij}t_{bhj}t_{cgi}$ are identically zero. Moreover, we have the tensor identity

$$(3.2) \quad t_{ea}(bt_{cd})_e = t_{e(ab}t_{cd})_e.$$

Here $X_{(a_1 \dots a_p)} := \sum_{\sigma \in S_p} \frac{1}{p!} X_{\sigma(a_1) \dots \sigma(a_p)}$.

Proof. By direct calculation in the defining matrix representation. \square

3.3. Quadratic and cubic states

With $N \in \mathbb{Z}_{\geq 1}$ as above, let $z_1, \dots, z_N \in \mathbb{C}$ be a collection of distinct points in the complex plane. For any $A \in \widehat{\mathfrak{sl}}_M$ let us define the $\widehat{\mathfrak{sl}}_M^{\oplus N}$ -valued rational function

$$(3.3) \quad A(z) := \sum_{i=1}^N \frac{A^{(i)}}{z - z_i}.$$

Lemma 3.2. *For any $A, B \in \widehat{\mathfrak{sl}}_M$,*

$$\begin{aligned} [A(z), B(w)] &= -\frac{[A, B](z) - [A, B](w)}{z - w}, & z \neq w, \\ [A(z), B(z)] &= -[A, B]'(z) \end{aligned}$$

where in the second line the $'$ denotes derivative with respect to the argument, z . \square

10 Sylvain Lacroix, Benoît Vicedo and Charles Young

Now we define states $\varsigma_1(z)$ and $\varsigma_2(z)$ in $\mathbb{V}_0^{\mathbf{k}}$, depending rationally on z , as follows

$$\begin{aligned}\varsigma_1(z) &:= \frac{1}{2} I_{-1}^a(z) I_{-1}^a(z) |0\rangle^{\mathbf{k}}, \\ \varsigma_2(z) &:= \frac{1}{3} t_{abc} I_{-1}^a(z) I_{-1}^b(z) I_{-1}^c(z) |0\rangle^{\mathbf{k}}.\end{aligned}$$

Define also the *twist function*

$$(3.4) \quad \varphi(z) := \sum_{i=1}^N \frac{k_i}{z - z_i}.$$

and, for any integer j , the *twisted derivative operator of degree j with respect to z* ,

$$(3.5) \quad D_z^{(j)} := \partial_z - \frac{j}{M} \varphi(z).$$

Recall the translation operator $T \in \text{End}(\mathbb{V}_0^{\mathbf{k}})$.

3.4. Main result

We can now state the main results of the paper.

Theorem 3.3. *For all $i, j \in \{1, 2\}$, there exist $\mathbb{V}_0^{\mathbf{k}}$ -valued rational functions $A_{ij}(z, w)$ and $B_{ij}(z, w)$ such that*

$$(3.6) \quad \varsigma_i(z)_{(0)} \varsigma_j(w) = (j D_z^{(i)} - i D_w^{(j)}) A_{ij}(z, w) + T B_{ij}(z, w).$$

Moreover, $A_{ij}(z, w)$ are regular at $z = w$ up to terms of the form TZ with $Z \in \mathbb{V}_0^{\mathbf{k}}$.

Proof. The first statement follows from a (lengthy) direct computation using the various identities from Lemma 3.1. Explicitly, we find that one choice of functions $A_{ij}(z, w)$ and $B_{ij}(z, w)$ is given by:

$$A_{11}(z, w) = \frac{M}{z - w} I_{-2}^a(z) I_{-1}^a(w) |0\rangle^{\mathbf{k}}, \quad B_{11}(z, w) = \frac{M}{(z - w)^2} I_{-1}^a(z) I_{-1}^a(w) |0\rangle^{\mathbf{k}}$$

Cubic hypergeometric integrals of motion in affine Gaudin models 11

for $i = j = 1$,

$$\begin{aligned} A_{12}(z, w) &= \frac{M}{2(z-w)} t_{abc} I_{-2}^a(z) I_{-1}^b(w) I_{-1}^c(w) |0\rangle^{\mathbf{k}}, \\ B_{12}(z, w) &= \frac{M}{2(z-w)^2} t_{abc} I_{-1}^a(z) I_{-1}^b(w) I_{-1}^c(w) |0\rangle^{\mathbf{k}}, \end{aligned}$$

for $i = 1$ and $j = 2$,

$$\begin{aligned} A_{21}(z, w) &= \frac{M}{z-w} t_{abc} I_{-2}^a(z) I_{-1}^b(z) I_{-1}^c(w) |0\rangle^{\mathbf{k}}, \\ B_{21}(z, w) &= \frac{M}{(z-w)^2} t_{abc} I_{-1}^a(z) I_{-1}^b(z) I_{-1}^c(w) |0\rangle^{\mathbf{k}}, \end{aligned}$$

for $i = 2$ and $j = 1$, and finally

$$\begin{aligned} A_{22}(z, w) &= \frac{M}{2(z-w)} t_{abe} t_{cde} I_{-2}^a(z) I_{-1}^b(z) I_{-1}^c(w) I_{-1}^d(w) |0\rangle^{\mathbf{k}} \\ &\quad - \frac{M(M^2-4)}{(z-w)^2} \left(\frac{\varphi(z) - \varphi(w)}{M} I_{-4}^a(z) I_{-1}^a(w) \right. \\ &\quad \left. + \frac{1}{2} (I_{-4}^a(z) I_{-1}^a(w) + I_{-4}^a(z) I_{-1}^a(w)) - \frac{1}{M} f_{abc} I_{-3}^a(z) I_{-1}^b(z) I_{-1}^c(w) \right) |0\rangle^{\mathbf{k}} \\ (3.7) \quad &\quad + \frac{M(M^2-4)}{(z-w)^3} (I_{-4}^a(z) I_{-1}^a(z) - 3I_{-4}^a(z) I_{-1}^a(w) - I_{-3}^a(z) I_{-2}^a(z)) |0\rangle^{\mathbf{k}} \end{aligned}$$

and

$$\begin{aligned} B_{22}(z, w) &= \frac{M}{2(z-w)^2} t_{abe} t_{cde} I_{-1}^a(z) I_{-1}^b(z) I_{-1}^c(w) I_{-1}^d(w) |0\rangle^{\mathbf{k}} \\ &\quad - \frac{2M(M^2-4)}{(z-w)^3} \left(2 \frac{\varphi(z) - \varphi(w)}{M} I_{-3}^a(z) I_{-1}^a(w) + I_{-3}^a(z) I_{-1}^a(w) \right. \\ &\quad \left. - \frac{1}{M} f_{abc} I_{-2}^a(z) I_{-1}^b(z) I_{-1}^c(w) \right) |0\rangle^{\mathbf{k}} \\ &\quad + \frac{2M(M^2-4)}{(z-w)^4} (I_{-3}^a(z) I_{-1}^a(z) - 5I_{-3}^a(z) I_{-1}^a(w) \\ &\quad - I_{-2}^a(z) I_{-2}^a(z) + \frac{1}{2} I_{-2}^a(z) I_{-2}^a(w)) |0\rangle^{\mathbf{k}}. \end{aligned}$$

in the case $i = j = 2$.

To show the ‘moreover’ part, it suffices to expand the expressions for $A_{ij}(z, w)$ given above in w near z and express all singular terms in the

12 Sylvain Lacroix, Benoît Vicedo and Charles Young

desired form. When $i = j = 1$ we have

$$A_{11}(z, w) = \frac{M}{2(z-w)} T I_{-1}^a(z) I_{-1}^a(z) |0\rangle^{\mathbf{k}} + \dots$$

where the dots represent terms regular at $z = w$. Likewise, for $i = 1, j = 2$ and $i = 2, j = 1$ we find, respectively,

$$\begin{aligned} A_{12}(z, w) &= \frac{M}{3(z-w)} t_{abc} T I_{-1}^a(z) I_{-1}^b(z) I_{-1}^c(z) |0\rangle^{\mathbf{k}} + \dots \\ A_{21}(z, w) &= \frac{M}{3(z-w)} t_{abc} T I_{-1}^a(z) I_{-1}^b(z) I_{-1}^c(z) |0\rangle^{\mathbf{k}} + \dots \end{aligned}$$

as required.

Now consider the case $i = j = 2$. The terms of order $(z-w)^{-3}$ and $(z-w)^{-2}$ in $A_{22}(z, w)$ can be written as

$$\begin{aligned} & - \frac{M(M^2-4)}{3(z-w)^3} T(2I_{-3}^a(z)I_{-1}^a(z) + \tfrac{1}{4}I_{-2}^a(z)I_{-2}^a(z)) |0\rangle^{\mathbf{k}} \\ & + \frac{M(M^2-4)}{6(z-w)^2} T(5I_{-3}^a(z)I_{-1}^a{}'(z) - I_{-3}^a{}'(z)I_{-1}^a(z) + \tfrac{1}{2}I_{-2}^a{}'(z)I_{-2}^a(z)) |0\rangle^{\mathbf{k}}. \end{aligned}$$

The terms of order $(z-w)^{-1}$ in $A_{22}(z, w)$ can also be written as TZ for some state $Z \in \mathbb{V}_0^{\mathbf{k}}$. In order to see this, one needs to note that the singular part of the first term on the right hand side of (3.7) can be written as

$$\begin{aligned} & \frac{M}{z-w} t_{abe} t_{cde} I_{-2}^a(z) I_{-1}^b(z) I_{-1}^c(z) I_{-1}^d(z) |0\rangle^{\mathbf{k}} \\ & = \frac{M}{z-w} t_{abe} t_{cde} I_{-2}^a(z) I_{-1}^{(b}(z) I_{-1}^c(z) I_{-1}^{d)}(z) |0\rangle^{\mathbf{k}} \\ & \quad - \frac{M(M^2-4)}{z-w} f_{abc} \left(\frac{2}{3} I_{-2}^a(z) I_{-2}^c{}'(z) I_{-1}^b(z) + \frac{1}{3} I_{-2}^a(z) I_{-1}^b(z) I_{-2}^c{}'(z) \right) |0\rangle^{\mathbf{k}}. \end{aligned}$$

The first term on the right hand side is proportional to

$$\begin{aligned} t_{abe} t_{cde} I_{-2}^a(z) I_{-1}^{(b}(z) I_{-1}^c(z) I_{-1}^{d)}(z) |0\rangle^{\mathbf{k}} &= t_{e(ab} t_{cd)e} I_{-2}^a(z) I_{-1}^b(z) I_{-1}^c(z) I_{-1}^d(z) |0\rangle^{\mathbf{k}} \\ &= \frac{1}{4} t_{e(ab} t_{cd)e} T I_{-1}^a(z) I_{-1}^b(z) I_{-1}^c(z) I_{-1}^d(z) |0\rangle^{\mathbf{k}} \end{aligned}$$

where in the first equality we used the identity (3.2). \square

Theorem 3.4. *For any $x \in \mathfrak{sl}_M$ and $n \in \mathbb{Z}_{\geq 0}$ we have*

$$\begin{aligned}\Delta x_n \varsigma_1(z) &= D_z^{(1)}(-Mx_{-1}(z)|0\rangle \delta_{n,1}) \\ \Delta x_n \varsigma_2(z) &= D_z^{(2)}(-\tfrac{1}{2}Mt_{abc}(x|I^a)I_{-1}^b(z)I_{-1}^c(z)|0\rangle \delta_{n,1}).\end{aligned}$$

Proof. As for Theorem 3.3, the proof is by direct calculation using the identities in Lemma 3.1. \square

3.5. Quadratic Gaudin Hamiltonians

Recall from §3.1 that we are assuming the levels to be non-critical, *i.e.* $k_i \neq -M$ for each $i \in \{1, \dots, N\}$. Let

$$(3.8) \quad \omega^{(i)} := \frac{1}{2(k_i + M)} I_{-1}^{a(i)} I_{-1}^{a(i)} |0\rangle^{\mathbf{k}} \in \mathbb{V}_0^{\mathbf{k}}$$

be the corresponding Segal-Sugawara state at site i . It has the property that for any $A = \bigotimes_{i=1}^N A_i \in \mathbb{V}_0^{\mathbf{k}}$ we have $(\omega^{(i)})_{(0)} A = A_1 \otimes \dots \otimes A_{i-1} \otimes T A_i \otimes A_{i+1} \otimes \dots \otimes A_N$. Let us also introduce the $\mathbb{V}_0^{\mathbf{k}}$ -valued rational function

$$\omega(z) := \sum_{i=1}^N \frac{\omega^{(i)}}{z - z_i}.$$

Consider the state

$$(3.9) \quad s_1(z) := \varsigma_1(z) + M D_z^{(1)} \omega(z) \in \mathbb{V}_0^{\mathbf{k}}$$

depending rationally on z .

Theorem 3.5. *For $i \in \{1, 2\}$, we have*

$$s_1(z)_{(0)} \varsigma_i(w) = -D_w^{(i)} A_{1i}(z, w) + T \left(B_{1i}(z, w) + D_z^{(1)} \frac{M \varsigma_i(w)}{z - w} \right),$$

with the $\mathbb{V}_0^{\mathbf{k}}$ -valued rational functions $A_{1i}(z, w)$ and $B_{1i}(z, w)$ as in Theorem (3.3).

14 Sylvain Lacroix, Benoît Vicedo and Charles Young

Proof. We have

$$\begin{aligned}\omega(z)_{(0)\varsigma_2}(w) &= -\frac{1}{z-w} t_{abc}(I_{-2}^a(z) - I_{-2}^a(w)) I_{-1}^b(w) I_{-1}^c(w) |0\rangle^k \\ &= -\frac{2}{M} A_{12}(z, w) + T\left(\frac{\varsigma_2(w)}{z-w}\right).\end{aligned}$$

Similarly, one finds that

$$\omega(z)_{(0)\varsigma_1}(w) = -\frac{1}{M} A_{11}(z, w) + T\left(\frac{\varsigma_1(w)}{z-w}\right).$$

Acting with the twisted derivative $D_z^{(1)}$ on the above equations we obtain

$$(MD_z^{(1)}\omega(z))_{(0)\varsigma_j}(w) = -jD_z^{(1)}A_{1j}(z, w) + TD_z^{(1)}\frac{M\varsigma_j(w)}{z-w},$$

for $j \in \{1, 2\}$. The result now follows from adding this to (3.6) with $i = 1$ and using the definition (3.9) of the state $s_1(z)$. \square

3.6. Twisted cycles

Let us define

$$\mathcal{P}(z) := \prod_{j=1}^N (z - z_j)^{k_j}.$$

It is a multivalued function on $\mathbb{C} \setminus \{z_1, \dots, z_N\}$.

Pick $i \in \mathbb{Z}$. Let γ be a closed contour in $\mathbb{C} \setminus \{z_1, \dots, z_N\}$ along which there exists a univalued branch of the function $\mathcal{P}(z)^{i/M}$. For example one can let γ be a Pochhammer contour about any pair of the marked points z_1, \dots, z_N . The pair, consisting of such a contour γ and univalued branch of $\mathcal{P}(z)^{i/M}$ along it, defines a cycle (possibly the zero cycle) in the twisted homology corresponding to $D_z^{(i)}$ – see [LVY] for the precise definition.

Lemma 3.6. *One has*

$$\int_{\gamma} \mathcal{P}(z)^{-i/M} D_z^{(i)} f(z) dz = 0$$

for any meromorphic f which is nonsingular on γ .

Proof. This follows by noting that $D_z^{(i)} f(z) = \mathcal{P}(z)^{i/M} \partial_z (\mathcal{P}(z)^{-i/M} f(z))$. \square

3.7. Quadratic and cubic Hamiltonians

Suppose $X, Y \in \mathbb{V}_0^k$ are any two states. Consider the formal modes $X_{(0)}, Y_{(0)} \in \tilde{U}_k(\widehat{\mathfrak{sl}}_M^{\oplus N})$. One can show that

$$[X_{(0)}, Y_{(0)}] = (X_{(0)}Y)_{(0)}.$$

This is actually a general statement which holds for any vertex algebra. See, for example, [FBZ04, Chapter 4].¹

It follows that, given a collection of states in \mathbb{V}_0^k whose zeroth products vanish, the formal zero modes of these states define a commutative subalgebra of the algebra of observables $\tilde{U}_k(\widehat{\mathfrak{sl}}_M^{\oplus N})$. In fact the zeroth products do not even need to vanish for this to be true: one can show that $(TZ)_{(0)} = 0$ for any state Z , so it is sufficient that the zeroth products are translates, *i.e.* lie in the image of T .

In our case we have the states $\varsigma_1(z)$ and $\varsigma_2(z)$ which depend rationally on the spectral parameter $z \in \mathbb{C} \setminus \{z_1, \dots, z_N\}$. Their zeroth products, according to Theorem 3.3, are zero modulo translates *and* twisted derivatives in the spectral parameters. Let γ_i and η_i be any cycles of the twisted homology corresponding to $D_z^{(i)}$, for $i \in \{1, 2\}$. In view of Lemma 3.6, we have

$$(3.10) \quad \left[\int_{\gamma_i} \mathcal{P}(z)^{-i/M} \varsigma_i(z)_{(0)} dz, \int_{\eta_j} \mathcal{P}(w)^{-j/M} \varsigma_j(w)_{(0)} dw \right] = 0$$

for $i, j \in \{1, 2\}$, and we get a commutative subalgebra of the algebra of observables $\tilde{U}_k(\widehat{\mathfrak{sl}}_M^{\oplus N})$, generated by such integrals over twisted cycles.

However, these are not quite the Hamiltonians we want, because their action on highest weight modules is not diagonalizable. Indeed, recall that the *homogeneous gradation* on $U(\widehat{\mathfrak{sl}}_M^{\oplus N})$ is the \mathbb{Z} -gradation in which $\deg(X_n^{(i)}) = n$. It induces a $\mathbb{Z}_{\leq 0}$ -gradation on \mathbb{V}_0^k , if we set $\deg(|0\rangle^k) = 0$. One can show that if $X \in \mathbb{V}_0^k$ has $\deg(X) = k$ then $\deg(X_{(n)}) = 1 + k + n$. Now in our case $\deg(\varsigma_i(z)) = -i - 1$, and hence

$$\deg(\varsigma_i(z)_{(0)}) = -i$$

for $i \in \{1, 2\}$. Thus the operators $\int_{\gamma_i} \mathcal{P}(z)^{-i/M} \varsigma_i(z)_{(0)} dz \in \tilde{U}_k(\widehat{\mathfrak{sl}}_M^{\oplus N})$ are of degree $-i \neq 0$ in the homogeneous gradation. Consequently if M is any

¹Note that there the notation for formal modes is *e.g.* $X_{[0]}$, with $X_{(0)}$ reserved for the representative in $\text{End } \mathbb{V}$.

16 Sylvain Lacroix, Benoît Vicedo and Charles Young

graded module over $\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ whose subspace of grade n is trivial for sufficiently large n then the operator $\int_{\gamma_i} \mathcal{P}(z)^{-i/M} \varsigma_i(z)_{(0)} dz$ has no non-zero generalised eigenvalue in M .

Instead, what we want are commuting operators in $\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ that are of degree zero in the homogeneous gradation. There is a general procedure to construct such operators, which roughly speaking involves a change of coordinate² to go from the plane to the cylinder [Nah91]. For our present purposes we just recall the following facts: there is a notion of what we shall call the *Fourier modes*, $X_n \in \tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$, of any state $X \in \mathbb{V}_0^{\mathbf{k}}$. These modes have the property that X_n always has grade n . They obey the commutator formula

$$(3.11) \quad [X_m, Y_n] = \left(X_{(0)}Y + \sum_{k=1}^{\infty} \frac{m^k X_{(k)}Y}{k!} \right)_{m+n}$$

for any states $X, Y \in \mathbb{V}_0^{\mathbf{k}}$ and any $n, m \in \mathbb{Z}$. In particular

$$[X_0, Y_0] = (X_{(0)}Y)_0.$$

Also, $(TZ)_0 = 0$ for any state $Z \in \mathbb{V}_0^{\mathbf{k}}$.

One has $(x_{-1}^{(i)} | 0 \rangle^{\mathbf{k}})_n = x_n$ for $x \in \mathfrak{sl}_M$ and there is a formula, analogous to the normal ordered product formula (2.4), which allows one to compute by recursion the Fourier modes of a general state $X \in \mathbb{V}_0^{\mathbf{k}}$. (See for example equation (7) in [BLZ96].) In general the resulting expressions for modes of states are rather intricate but for $\varsigma_1(z)_0$ and $\varsigma_2(z)_0$ they are simple: one finds

$$(3.12) \quad \begin{aligned} \varsigma_1(z)_0 &= \frac{\dim(\mathfrak{sl}_M)}{24} \varphi'(z) + \frac{1}{2} I_0^a(z) I_0^a(z) + \sum_{n>0} I_{-n}^a(z) I_n^a(z), \\ \varsigma_2(z)_0 &= \frac{1}{3} t_{abc} \sum_{j,k \geq 0} \left(I_{-1-k}^a(z) I_{-1-j}^b(z) I_{2+j+k}^c(z) + 2 I_{-1-k}^a(z) I_{1+k-j}^b(z) I_j^c(z) \right. \\ (3.13) \quad &\quad \left. + I_{-j-k}^a(z) I_k^b(z) I_j^c(z) \right). \end{aligned}$$

Let us now define

$$(3.14) \quad \hat{Q}_i^\gamma := \int_{\gamma} \mathcal{P}(z)^{-i/M} \varsigma_i(z)_0 dz$$

²(in the “loop variable”, t , not the “spectral variable” z)

Cubic hypergeometric integrals of motion in affine Gaudin models 17

for $i \in \{1, 2\}$ and for γ any cycle of the twisted homology corresponding to $D_z^{(i)}$.

Recall the Lie algebra $\mathfrak{g} := \widehat{\mathfrak{sl}}_M \oplus \mathbb{C}d$ from §2.1. Let $\tilde{U}(\mathfrak{g}^{\oplus N}) := \varprojlim U(\mathfrak{g}^{\oplus N})/J_n^N$ be the completion of $U(\mathfrak{g}^{\oplus N})$ where for each $n \geq 0$, J_n^N is the ideal of $U(\mathfrak{g}^{\oplus N})$ generated by $a_r^{(i)}$ for all $r \geq n$, $a \in \mathfrak{sl}_M$ and $i \in \{1, \dots, N\}$. The quadratic Gaudin Hamiltonians are the elements of $\tilde{U}(\mathfrak{g}^{\oplus N})$ defined as

$$\mathcal{H}_i := \sum_{\substack{j=1 \\ j \neq i}}^N \frac{k^{(i)}d^{(j)} + d^{(i)}k^{(j)} + \sum_{n \in \mathbb{Z}} I_{-n}^{a(i)} I_n^{a(j)}}{z_i - z_j}, \quad i = 1, \dots, N.$$

For each $i \in \{1, \dots, N\}$ we also have the i^{th} copy of the quadratic Casimir of \mathfrak{g} in $\tilde{U}(\mathfrak{g}^{\oplus N})$, which is defined as

$$\mathcal{C}^{(i)} := (k^{(i)} + M)d^{(i)} + \frac{1}{2}I_0^{a(i)}I_0^{a(i)} + \sum_{n>0} I_{-n}^{a(i)}I_n^{a(i)}.$$

The algebra $\tilde{U}_k(\widehat{\mathfrak{sl}}_M^{\oplus N})$, introduced in §3.1, is isomorphic to the quotient of $\tilde{U}(\mathfrak{g}^{\oplus N})$ by the ideal generated by $k^{(i)} - k_i$ and $\mathcal{C}^{(i)}$ for each $i \in \{1, \dots, N\}$.

Recall the state $s_1(z) \in \mathbb{V}_0^k$ defined in (3.9).

Lemma 3.7. *The operator $s_1(z)_0$ is the image in $\tilde{U}_k(\widehat{\mathfrak{sl}}_M^{\oplus N})$ of the operator*

$$\sum_{i=1}^N \frac{\mathcal{C}^{(i)}}{(z - z_i)^2} + \sum_{i=1}^N \frac{\mathcal{H}_i}{z - z_i} \in \tilde{U}(\mathfrak{g}^{\oplus N}).$$

Proof. Computing the Fourier zero mode of the state $s_1(z)$ we find

$$\begin{aligned} s_1(z)_0 &= \sum_{i=1}^N \frac{1}{(z - z_i)^2} \left(-\frac{\dim(\mathfrak{sl}_M)}{24} k_i + \frac{1}{2} I_0^{a(i)} I_0^{a(i)} + \sum_{n>0} I_{-n}^{a(i)} I_n^{a(i)} \right) \\ &\quad + \sum_{i=1}^N \frac{1}{z - z_i} \sum_{\substack{j=1 \\ j \neq i}}^N \frac{1}{z_i - z_j} \sum_{n \in \mathbb{Z}} I_{-n}^{a(i)} I_n^{a(j)}. \end{aligned}$$

On the other hand, we also have

$$MD_z^{(1)} \omega(z)_0 = - \sum_{i=1}^N \frac{(k_i + M)(\omega^{(i)})_0}{(z - z_i)^2} - \sum_{i=1}^N \frac{1}{z - z_i} \sum_{\substack{j=1 \\ j \neq i}}^N \frac{k_i(\omega^{(j)})_0 + k_j(\omega^{(i)})_0}{z_i - z_j}.$$

18 Sylvain Lacroix, Benoît Vicedo and Charles Young

The result now follows from combining the above and noting that the Fourier zero mode of the Segal-Sugawara state $\omega^{(i)}$, for each $i \in \{1, \dots, N\}$, as defined in (3.8), reads

$$(\omega^{(i)})_0 = \frac{\frac{1}{2}I_0^{a(i)}I_0^{a(i)} + \sum_{n>0}I_{-n}^{a(i)}I_n^{a(i)}}{k_i + M} - \frac{k_i \dim(\mathfrak{sl}_M)}{24(k_i + M)}.$$

□

We now have the following corollary of Theorems 3.3, 3.4 and 3.5.

Theorem 3.8. *The operators $\hat{Q}_i^\gamma \in \tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ defined in (3.14) are of grade 0 in the homogeneous gradation. They have the following properties:*

- (i) *For $i, j \in \{1, 2\}$ and any cycles γ and η of the twisted homology corresponding to $D_z^{(i)}$ and $D_z^{(j)}$, respectively, we have*

$$[\hat{Q}_i^\gamma, \hat{Q}_j^\eta] = 0.$$

- (ii) *For all $i \in \{1, 2\}$, any cycle γ of the twisted homology corresponding to $D_z^{(i)}$ and $j \in \{1, \dots, N\}$ we have*

$$[\mathcal{H}_j, \hat{Q}_i^\gamma] = 0.$$

- (iii) *For all $i \in \{1, 2\}$, any cycle γ of the twisted homology corresponding to $D_z^{(i)}$ and all $x \in \mathfrak{sl}_M$ and $n \in \mathbb{Z}$, we have*

$$[\Delta x_n, \hat{Q}_i^\gamma] = 0.$$

□

According to the conjecture in [LVY], the Hamiltonians \hat{Q}_i^γ can be diagonalized by means of the Bethe ansatz and have certain specific eigenvalues encoded by affine opers. For \hat{Q}_1^γ this is known; see [LVY, §5]. In the next two sections we perform a simple Bethe ansatz calculation to check that the conjecture is correct also for \hat{Q}_2^γ at least for zero and one Bethe roots.

4. Opers and eigenvalues

The general expectation is that, following the situation in finite types – for which see *e.g.* [Fre05] – the spectrum of Hamiltonians for the Gaudin model associated with \mathfrak{g} should be described in terms of *opers* for the Langlands dual Lie algebra ${}^L\mathfrak{g}$, *i.e.* the Lie algebra with transposed Cartan matrix. Since

Cubic hypergeometric integrals of motion in affine Gaudin models 19

$\widehat{\mathfrak{sl}}_M$ has symmetric Cartan matrix it is self-dual: ${}^L\widehat{\mathfrak{sl}}_M \cong \widehat{\mathfrak{sl}}_M$. Nevertheless, to keep the general structure in sight and to follow the notation of [LVY] we prefer to distinguish the two copies of $\widehat{\mathfrak{sl}}_M$. Thus, let \check{e}_i, \check{f}_i be the Chevalley-Serre generators of the dual copy ${}^L\widehat{\mathfrak{sl}}_M$. They obey the same relations (2.2) as the generators e_i, f_i above but with the roots and coroots α_i and $\check{\alpha}_i$ interchanged. The Cartan subalgebra of ${}^L\widehat{\mathfrak{sl}}_M$ is ${}^L\mathfrak{h} := \mathfrak{h}^*$. Let ${}^L\widehat{\mathfrak{sl}}_M = {}^L\mathfrak{n}_- \oplus {}^L\mathfrak{h} \oplus {}^L\mathfrak{n}_+$ be the Cartan decomposition.

The paper [LVY] contains an explicit conjecture of how ${}^L\widehat{\mathfrak{sl}}_M$ -opers encode the eigenvalues of the Hamiltonians above, as we now recall.

4.1. Opers and Miura opers of type ${}^L\widehat{\mathfrak{sl}}_M$

Pick a derivation element $\rho \in {}^L\mathfrak{h}$ such that $[\rho, \check{e}_i] = \check{e}_i, [\rho, \check{f}_i] = -\check{f}_i$ for each $i = 0, 1, \dots, M-1$. Explicitly, one has

$$(4.1) \quad \rho := Md + \dot{\rho}$$

where $\dot{\rho} := \sum_{i=1}^{M-1} \omega_i \in \mathfrak{h}^*$ is the Weyl vector of \mathfrak{sl}_M .

Let $p_{-1} := \sum_{i=0}^{M-1} \check{f}_i \in {}^L\mathfrak{n}_-$. A *Miura oper* of type ${}^L\widehat{\mathfrak{sl}}_M$ is a connection of the form

$$(4.2a) \quad d + (p_{-1} + u(z)) dz$$

where $u(z)$ is a meromorphic function valued in ${}^L\mathfrak{h} = \mathfrak{h}^*$. Since $\{\alpha_i\}_{i=0}^{M-1} \cup \{\rho\}$ is a basis of ${}^L\mathfrak{h}$ we have

$$(4.2b) \quad u(z) = \sum_{i=0}^{M-1} u_i(z) \alpha_i - \frac{\varphi(z)}{M} \rho$$

for some meromorphic functions $\{u_i(z)\}_{i=0}^{M-1}$ and $\varphi(z)$. An ${}^L\widehat{\mathfrak{sl}}_M$ -oper is a gauge equivalence class of connections of the form

$$d + (p_{-1} + b(z)) dz,$$

where $b(z)$ is a meromorphic function valued in ${}^L\mathfrak{b}_+ = {}^L\mathfrak{h} \oplus {}^L\mathfrak{n}_+$, under the gauge action of the group ${}^L N_+ = \exp({}^L\mathfrak{n}_+)$. Any ${}^L\widehat{\mathfrak{sl}}_M$ -oper has a *quasi-canonical* representative [LVY]: namely, by suitable gauge transformation

20 Sylvain Lacroix, Benoît Vicedo and Charles Young

the connection can be brought to the form

$$(4.3) \quad \nabla = d + \left(p_{-1} - \frac{\varphi}{M} \rho + \sum_{j \in E} v_j p_j \right) dz,$$

where the sum is over the set

$$E = \mathbb{Z}_{\geq 1} \setminus M\mathbb{Z}_{\geq 1}$$

of positive *exponents* of ${}^L\widehat{\mathfrak{sl}}_M$, and where for each $j \in E$, v_j is a meromorphic function. For each $j \in \pm E$, p_j is a certain element of ${}^L\mathfrak{n}_+$ of grade j in the principal gradation, *i.e.*

$$[\rho, p_j] = j p_j,$$

such that the following commutation relations hold with p_{-1} :

$$[p_{-1}, p_j] = \begin{cases} 0 & j \in \pm E \setminus \{1\} \\ -\delta & j = 1, \end{cases}$$

where $\delta := \sum_{i=0}^{M-1} \alpha_i$ is central in ${}^L\widehat{\mathfrak{sl}}_M$. It can be shown that, for each $j \in \pm E$, a non-zero such p_j exists and is unique up to normalization. Let $\mathfrak{a} := \mathbb{C}\{p_j\}_{j \in \pm E} \oplus \mathbb{C}\delta \oplus \mathbb{C}\rho \subset {}^L\widehat{\mathfrak{sl}}_M$. Then \mathfrak{a} is a Lie subalgebra of ${}^L\widehat{\mathfrak{sl}}_M$ called the *principal subalgebra* (the remaining non-trivial commutation relations are $[p_m, p_n] = m\delta_{m+n,0}\delta$). The normalization of the p_j can be so chosen that the restriction to \mathfrak{a} of the bilinear form $(\cdot|\cdot)$ is given by

$$(\delta|\rho) = (\rho|\delta) = M, \quad (p_m|p_n) = M\delta_{m+n,0}, \quad m, n \in \pm E.$$

The expression for v_1 below can be found in [LVY, Proposition 3.10].

Proposition 4.1. *The coefficients of p_1 and p_2 of any quasi-canonical form of the ${}^L\widehat{\mathfrak{sl}}_M$ -oper defined by the Miura ${}^L\widehat{\mathfrak{sl}}_M$ -oper (4.2) are given by*

$$\begin{aligned} v_1 &= \frac{1}{M} \left(\frac{1}{2}(u|u) + D_z^{(1)}(\rho|u) \right), \\ v_2 &= \frac{1}{M} \left(- \sum_{i=0}^{M-1} u_i(u_{i+1}^2 - u_{i-1}^2) - \frac{1}{2} \sum_{i=0}^{M-1} u_i(u'_{i+1} - u'_{i-1}) + D_z^{(2)}f_2 \right), \end{aligned}$$

where f_2 is an arbitrary meromorphic function.

Cubic hypergeometric integrals of motion in affine Gaudin models 21

Proof. Let $g(z) = \exp(m(z)) \in {}^L N_+$ with m a meromorphic function valued in ${}^L \mathfrak{n}_+$. Let $m = \sum_{n=1}^{\infty} m_n$ be the decomposition of m in the principal gradation.

Let $\nabla_{\partial_z} = \partial_z + p_{-1} + u - \frac{\varphi}{M}\rho$, as in (4.2). We demand that

$$(4.4) \quad g \nabla_{\partial_z} g^{-1} = \partial_z + p_{-1} - \frac{\varphi}{M}\rho + v_1 p_1 + v_2 p_2 + \dots$$

where the dots denote terms of degree at least three in the principal gradation. In grade 0 we have $u + [m_1, p_{-1}] - \frac{\varphi}{M}\rho = -\frac{\varphi}{M}\rho$ and hence $u = -[m_1, p_{-1}]$. Yet since $u = \sum_{i=0}^{M-1} u_i \alpha_i$, we must set

$$m_1 = - \sum_{i=0}^{M-1} u_i \check{e}_i.$$

Now consider the component of equation (4.4) in grade 1. We need

$$(4.5) \quad [m_2, p_{-1}] + \frac{1}{2}[m_1, [m_1, p_{-1}]] + [m_1, u] + \frac{\varphi}{M}m_1 - m'_1 = v_1 p_1.$$

Consider taking the overlap $(p_{-1}|\cdot)$ of both sides of this equation. Recall the definition of $D_z^{(j)}$ from (3.5). Since $(p_{-1}|p_1) = M$ we find, using the invariance of the bilinear form, that

$$\begin{aligned} Mv_1 &= -\frac{1}{2}([m_1, p_{-1}]|[m_1, p_{-1}]) - ([m_1, p_{-1}]|u) - D_z^{(1)}(p_{-1}|m_1) \\ &= \frac{1}{2}(u|u) + D_z^{(1)}(\rho|u), \end{aligned}$$

where $(p_{-1}|m_1) = ([p_{-1}, \rho]|m_1) = (\rho|[m_1, p_{-1}]) = -(\rho|u)$. (The equation (4.5) also fixes the component of m_2 in the orthogonal complement \mathfrak{c}_2 of \mathfrak{a}_2 , but we won't need that here.)

We turn to the component of equation (4.4) in grade 2. We have

$$\begin{aligned} (g \partial_z g^{-1})_2 &= -m'_2 - \frac{1}{2}[m_1, m'_1], \\ (gp_{-1}g^{-1})_2 &= [m_3, p_{-1}] + \frac{1}{2}[m_1, [m_2, p_{-1}]] + \frac{1}{2}[m_2, [m_1, p_{-1}]] + \frac{1}{6}[m_1, [m_1, [m_1, p_{-1}]]] \\ &= [m_3, p_{-1}] + \frac{1}{2}[[m_1, m_2], p_{-1}] - [m_2, u] - \frac{1}{6}[m_1, [m_1, u]], \\ (gug^{-1})_2 &= [m_2, u] + \frac{1}{2}[m_1, [m_1, u]], \\ -(g\rho g^{-1})_2 &= -[m_2, \rho] - \frac{1}{2}[m_1, [m_1, \rho]] = 2m_2. \end{aligned}$$

Hence we need

$$[m_3 + \frac{1}{2}[m_1, m_2], p_{-1}] + \frac{1}{3}[m_1, [m_1, u]] - \frac{1}{2}[m_1, m'_1] - D_z^{(2)}m_2 = v_2 p_2.$$

22 Sylvain Lacroix, Benoît Vicedo and Charles Young

Taking $(p_{-2}|\cdot)$ of both sides and then using $(p_{-2}|p_2) = M$, the invariance of $(\cdot|\cdot)$ and the relation $[p_{-1}, p_{-2}] = 0$, we find

$$Mv_2 = -\frac{1}{3}([m_1, p_{-2}][m_1, u]) + \frac{1}{2}([m_1, p_{-2}]m'_1) - D_z^{(2)}(p_{-2}|m_2).$$

Now in fact, in type A , we have $p_{-2} = -\sum_{i=0}^{M-1}[\check{f}_i, \check{f}_{i+1}]$ (addition of indices modulo M). Hence

$$[m_1, p_{-2}] = \sum_{i,j=0}^{M-1} u_i[\check{e}_i, [\check{f}_j, \check{f}_{j+1}]] = -\sum_{i=0}^{M-1} u_i(\check{f}_{i-1} - \check{f}_{i+1})$$

and

$$[m_1, u] = -\sum_{i,j=0}^{M-1} u_i u_j[\check{e}_i, \alpha_j] = \sum_{i,j=0}^{M-1} a_{ij} u_i u_j \check{e}_i.$$

Since the coefficient of p_2 of an ${}^L\widehat{\mathfrak{sl}}_M$ -oper in quasi-canonical form is defined only up to twisted derivative, we therefore have

$$\begin{aligned} Mv_2 &= \frac{1}{3} \sum_{i,j,k=0}^{M-1} a_{ij}(\delta_{i,k-1} - \delta_{i,k+1}) u_i u_j u_k - \frac{1}{2} \sum_{i,j=0}^{M-1} (\delta_{i,j-1} - \delta_{i,j+1}) u_i u'_j + D_z^{(2)} f_2 \\ &= -\sum_{i=0}^{M-1} u_i(u_{i+1}^2 - u_{i-1}^2) - \frac{1}{2} \sum_{i=0}^{M-1} u_i(u'_{i+1} - u'_{i-1}) + D_z^{(2)} f_2, \end{aligned}$$

for an arbitrary meromorphic function f_2 . Here we used the fact that the Chevalley-Serre generators of ${}^L\widehat{\mathfrak{sl}}_M$ are normalised as $(\check{e}_i|\check{f}_j) = \delta_{ij}$. \square

5. Bethe Ansatz for 0 and 1 roots

5.1. Homogeneous vs. principal gradations

Let us write $q(z) \in \tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ for the Fourier zero mode of the state $\varsigma_2(z)$ as in (3.13),

$$\begin{aligned} q(z) := \varsigma_2(z)_0 &= \frac{1}{3} t_{abc} \sum_{j,k \geq 0} \left(I_{-1-k}^a(z) I_{-1-j}^b(z) I_{2+j+k}^c(z) \right. \\ &\quad \left. + 2I_{-1-k}^a(z) I_{1+k-j}^b(z) I_j^c(z) + I_{-j-k}^a(z) I_k^b(z) I_j^c(z) \right). \end{aligned}$$

Cubic hypergeometric integrals of motion in affine Gaudin models 23

By definition of $\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$, $q(z)$ is an infinite sum

$$q(z) = q(z)_0^{\text{hom}} + q(z)_1^{\text{hom}} + q(z)_2^{\text{hom}} + \dots$$

such that, for each $n \geq 1$ we have $q(z) = q(z)_0^{\text{hom}} + \dots + q(z)_{n-1}^{\text{hom}}$ modulo J_n^N . Recall that J_n^N is the left ideal in $U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ generated by elements of grade at least n in the homogeneous gradation. Let $J_n^{N, \text{prin}}$ denote the left ideal generated by elements of grade at least n in the principal gradation. For each $n \geq 1$, there is a nesting of ideals $J_{Mn}^{N, \text{prin}} \subset J_n^N \subset J_{Mn-M+1}^{N, \text{prin}}$. So we could equivalently have defined $\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$ as the completion with respect to the inverse system $U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})/J_n^{N, \text{prin}}$.

Let $\mathfrak{b}_- := \mathfrak{n}_- \oplus \mathfrak{h}$. Denote by $U_{\mathbf{k}}(\mathfrak{b}_-^{\oplus N})$ the quotient of $U(\mathfrak{b}_-^{\oplus N})$ by the two-sided ideal generated by $\mathbf{k}^{(i)} - k_i$ for all $i \in \{1, \dots, N\}$, cf. §3.1. We have the vector space isomorphism $U_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N}) \cong_{\mathbb{C}} U_{\mathbf{k}}(\mathfrak{b}_-^{\oplus N}) \otimes U(\mathfrak{n}_+^{\oplus N})$. Let $U(\mathfrak{n}_+^{\oplus N})_n$ denote the subspace of $U(\mathfrak{n}_+^{\oplus N})$ spanned by elements of total degree n in the principal gradation. We have the unique decomposition

$$(5.1) \quad q(z) = q(z)_0 + q(z)_1 + q(z)_2 + \dots$$

where $q(z)_n \in U_{\mathbf{k}}(\mathfrak{b}_-^{\oplus N}) \otimes U(\mathfrak{n}_+^{\oplus N})_n$.

Recall the definition (4.1) of the derivation element ρ which measures the grade in the principal gradation. Let $\check{\eta}_i \in \mathfrak{h}$ be the set of fundamental coweights adapted to the principal gradation; namely, such that $\{\check{\eta}_i\}_{i=0}^{M-1} \cup \{\mathbf{k}\} \subset \mathfrak{h}$ is the dual basis to $\{\alpha_i\}_{i=0}^{M-1} \cup \{\rho/M\} \subset \mathfrak{h}^*$:

$$\langle \alpha_i, \check{\eta}_j \rangle = \delta_{ij}, \quad \langle \rho, \check{\eta}_j \rangle = 0, \quad \langle \alpha_i, \mathbf{k} \rangle = 0, \quad \langle \rho, \mathbf{k} \rangle = M$$

for all $i, j \in \{0, 1, \dots, M-1\}$.

Proposition 5.1. *The first two terms in the sum (5.1) are*

$$(5.2) \quad \begin{aligned} q(z)_0 &= - \sum_{i,j,k=0}^{M-1} \check{\eta}_i(z) (\check{\eta}_{i+1}(z)^2 - \check{\eta}_{i-1}(z)^2) + \frac{1}{2} \sum_{i,j=0}^{M-1} \check{\eta}_i(z) (\check{\eta}'_{i+1}(z) - \check{\eta}'_{i-1}(z)), \\ q(z)_1 &= \sum_{i,j=0}^{M-1} f_i(z) (\check{\eta}_{i+1}(z) - \check{\eta}_{i-1}(z)) e_i(z), \end{aligned}$$

modulo twisted derivatives of degree 2.

Proof. By direct calculation. □

24 Sylvain Lacroix, Benoît Vicedo and Charles Young

5.2. Spin chain

For each $i \in \{1, \dots, N\}$ let $\lambda_i \in \mathfrak{h}^*$ be a weight of $\widehat{\mathfrak{sl}}_M$ such that

$$(5.3) \quad \langle \lambda_i, \mathbf{k} \rangle = k_i.$$

Let \mathbb{M}_{λ_i} denote the Verma module over \mathfrak{g} of highest weight λ_i and define

$$\mathbb{M}_{\boldsymbol{\lambda}} := \mathbb{M}_{\lambda_1} \otimes \dots \otimes \mathbb{M}_{\lambda_N}.$$

It is a smooth module over $\widehat{\mathfrak{sl}}_M^{\oplus N}$ of level \mathbf{k} , and therefore a module over the algebra of observables $\tilde{U}_{\mathbf{k}}(\widehat{\mathfrak{sl}}_M^{\oplus N})$. Let $|\boldsymbol{\lambda}\rangle$ denote the highest weight vector in $\mathbb{M}_{\boldsymbol{\lambda}}$.

5.3. Vacuum eigenvalues

For this subsection, consider the Miura ${}^L\widehat{\mathfrak{sl}}_M$ -oper (4.2), with

$$u(z) = u^{(0)}(z) := - \sum_{j=1}^N \frac{\lambda_j}{z - z_j}.$$

That is, $u_i(z) = - \sum_{j=1}^N \frac{\langle \lambda_j, \tilde{\eta}_i \rangle}{z - z_j}$ for each $i \in \{0, 1, \dots, M-1\}$, and $\varphi(z)$ is the twist function as in (3.4), by virtue of (5.3). Let $v_1(z)$ and $v_2(z)$ be the coefficients of a quasi-canonical form of the underlying ${}^L\widehat{\mathfrak{sl}}_M$ -oper, as in Proposition 4.1.

Proposition 5.2. *Up to twisted derivatives of degree 2, the vacuum eigenvalue of $q(z)$ is $-Mv_2(z)$, i.e.*

$$q(z) |\boldsymbol{\lambda}\rangle = -Mv_2(z) |\boldsymbol{\lambda}\rangle + D_z^{(2)} |\varepsilon(z)\rangle$$

for some vector $|\varepsilon(z)\rangle \in \mathbb{M}_{\boldsymbol{\lambda}}$ depending rationally on z . In particular, for any cycle γ in the twisted homology corresponding to $D_z^{(2)}$ we have

$$\hat{Q}_2^\gamma |\boldsymbol{\lambda}\rangle = -M \int_\gamma \mathcal{P}(z)^{-2/M} v_2(z) dz |\boldsymbol{\lambda}\rangle.$$

Cubic hypergeometric integrals of motion in affine Gaudin models 25

Proof. We have $q(z)_n |\lambda\rangle = 0$ for all $n \geq 1$, i.e. $q(z) |\lambda\rangle = q(z)_0 |\lambda\rangle$. By definition of the $\check{\eta}_i$,

$$\check{\eta}_i(z) |\lambda\rangle = -\langle u(z), \check{\eta}_i \rangle |\lambda\rangle = -u_i(z) |\lambda\rangle$$

and then the statement follows from Propositions 4.1 and 5.1. \square

5.4. Eigenvalues at depth 1

Let $w \in \mathbb{C} \setminus \{z_1, \dots, z_N\}$ be an additional point in the complex plane, distinct from the marked points z_1, \dots, z_N . Pick $n \in \{0, 1, \dots, M-1\}$. Now, for this subsection, we consider the Miura ${}^L\widehat{\mathfrak{sl}}_M$ -oper, (4.2), with

$$(5.4) \quad u(z) = u^{(0)}(z) + \frac{\alpha_n}{z-w} = -\sum_{j=1}^N \frac{\lambda_j}{z-z_j} + \frac{\alpha_n}{z-w}$$

and let $v_1(z)$ and $v_2(z)$ be the coefficients of a quasi-canonical form of its underlying ${}^L\widehat{\mathfrak{sl}}_M$ -oper, as in Proposition 4.1.

The *weight function* or *Schechtman-Varchenko vector* for one Bethe root of colour $n \in \{0, 1, \dots, M-1\}$ is by definition $f_n(w) |\lambda\rangle$, where

$$f_n(w) = \sum_{j=1}^N \frac{f_n^{(i)}}{w-z_j}$$

using the notation (3.3).

Proposition 5.3. *Suppose w obeys the Bethe equation*

$$0 = \sum_{j=1}^N \frac{\langle \lambda_j, \check{\alpha}_n \rangle}{w-z_j}.$$

Then, up to twisted derivatives of degree 2, the weight function $f_n(w) |\lambda\rangle$ is an eigenstate of $q(z)$ with eigenvalue $-Mv_2(z)$, i.e.

$$q(z)f_n(w) |\lambda\rangle = -Mv_2(z)f_n(w) |\lambda\rangle + D_z^{(2)}|\tilde{\varepsilon}(z)\rangle,$$

for some vector $|\tilde{\varepsilon}(z)\rangle \in \mathbb{M}_\lambda$ depending rationally on z . In particular, for any cycle γ in the twisted homology corresponding to $D_z^{(2)}$ we have

$$\hat{Q}_2^\gamma f_n(w) |\lambda\rangle = -M \int_\gamma \mathcal{P}(z)^{-2/M} v_2(z) dz f_n(w) |\lambda\rangle.$$

26 Sylvain Lacroix, Benoît Vicedo and Charles Young

Proof. By definition $[\check{\eta}_i, f_n] = -f_n \langle \alpha_n, \check{\eta}_i \rangle = -\delta_{in} f_n$. Using this and Lemma 3.2, one computes, starting from the expressions for $q(z)_0$ in Proposition 5.1,

$$\begin{aligned} [q(z)_0, f_n(w)] |\lambda\rangle &= \left(- \sum_{i=0}^{M-1} a_{ni} u_i^{(0)}(z) \frac{u_{n+1}^{(0)}(z) - u_{n-1}^{(0)}(z)}{z - w} \right. \\ &\quad \left. - \frac{u_{n+1}^{(0)}(z) - u_{n-1}^{(0)}(z)}{2(z - w)^2} + \frac{u_{n+1}^{(0)'}(z) - u_{n-1}^{(0)'}(z)}{2(z - w)} \right) f_n(w) |\lambda\rangle \\ &\quad + \left(\sum_{i=0}^{M-1} a_{ni} u_i^{(0)}(z) \frac{u_{n+1}^{(0)}(z) - u_{n-1}^{(0)}(z)}{z - w} - \partial_z \frac{u_{n+1}^{(0)}(z) - u_{n-1}^{(0)}(z)}{z - w} \right) f_n(z) |\lambda\rangle. \end{aligned}$$

Here one recognizes the first term on the right hand side as the required correction to the vacuum value of $-Mv_2(z)$, cf. (5.4) and Proposition 4.1. Next, one finds

$$\begin{aligned} [q(z)_1, f_n(w)] |\lambda\rangle &= \sum_{i=0}^{M-1} [f_i(z) (\check{\eta}_{i+1}(z) - \check{\eta}_{i-1}(z)) e_i(z), f_n(w)] |\lambda\rangle \\ &= -f_n(z) (\check{\eta}_{n+1}(z) - \check{\eta}_{n-1}(z)) \frac{\check{\alpha}_n(z) - \check{\alpha}_n(w)}{z - w} |\lambda\rangle \\ &= -f_n(z) |\lambda\rangle (u_{n+1}^{(0)}(z) - u_{n-1}^{(0)}(z)) \frac{\langle u^{(0)}(z), \check{\alpha}_n \rangle - \langle u^{(0)}(w), \check{\alpha}_n \rangle}{z - w}. \end{aligned}$$

Now we have

$$\langle u(z)^{(0)}, \check{\alpha}_n \rangle = -\frac{\varphi(z)}{M} \langle \rho, \check{\alpha}_n \rangle + \sum_{i=0}^{M-1} \langle \alpha_i, \check{\alpha}_n \rangle u_i^{(0)}(z) = -\frac{\varphi(z)}{M} + \sum_{i=0}^{M-1} a_{ni} u_i^{(0)}(z).$$

Hence, one arrives at

$$\begin{aligned} q(z) f_n(w) |\lambda\rangle &= -Mv_2(z) f_n(w) |\lambda\rangle - D_z^{(2)} \left(\frac{1}{2} (u_{n+1}^{(0)}(z) - u_{n-1}^{(0)}(z)) \frac{f_n(z)}{z - w} |\lambda\rangle \right) \\ &\quad + \langle u^{(0)}(w), \check{\alpha}_n \rangle (u_{n+1}^{(0)}(z) - u_{n-1}^{(0)}(z)) \frac{f_n(z)}{z - w} |\lambda\rangle. \end{aligned}$$

The second line on the right hand side vanishes by the Bethe equation. \square

6. Two-point case, coset construction and the W3 algebra

In this section we specialize to the case of $N = 2$ marked points. For convenience, let us choose them to be $z_1 = 0$ and $z_2 = 1$. Let γ denote a Pochhammer contour around these two points. For $a \in \mathbb{C}$ and $n \in \mathbb{Z}_{\geq 0}$ let $(a)_n$ denote

Cubic hypergeometric integrals of motion in affine Gaudin models 27

the falling factorial

$$(a)_n := \prod_{k=0}^{n-1} (a - k) = a(a-1) \dots (a-n+1)$$

with $(a)_0 := 1$. For $a, b \in \mathbb{C}$, define

$$B(a, b) := \int_{\gamma} z^a (z-1)^b dz.$$

Lemma 6.1.

$$B(a-1, b) = \frac{a+b+1}{a} B(a, b), \quad B(a, b-1) = -\frac{a+b+1}{b} B(a, b).$$

Proof. Since the contour γ has no boundary, $aB(a-1, b) + bB(a, b-1) = 0$. From the identity $\frac{1}{z(z-1)} - \frac{1}{z-1} + \frac{1}{z} = 0$, one sees that $B(a-1, b-1) - B(a, b-1) + B(a-1, b) = 0$. The result follows. \square

Define

$$\begin{aligned} \omega &:= \omega^{(1)} + \omega^{(2)} - \omega^{(diag)} \\ &= \frac{1}{2(k_1 + M)} I_{-1}^{a(1)} I_{-1}^{a(1)} |0\rangle^{\mathbf{k}} + \frac{1}{2(k_1 + M)} I_{-1}^{a(2)} I_{-1}^{a(2)} |0\rangle^{\mathbf{k}} \\ &\quad - \frac{1}{2(k_1 + k_2 + M)} (I_{-1}^{a(1)} + I_{-1}^{a(2)}) (I_{-1}^{a(1)} + I_{-1}^{a(2)}) |0\rangle^{\mathbf{k}}. \end{aligned}$$

Define also the state $W \in \mathbb{V}_0^{\mathbf{k}}$ by

$$\begin{aligned} W &:= C(k_1, k_2) \left(\frac{1}{3} t_{abc} I_{-1}^{a(1)} I_{-1}^{b(1)} I_{-1}^{c(1)} |0\rangle^{\mathbf{k}} \left(-\frac{2}{M} k_2 \right) \left(-\frac{2}{M} k_2 - 1 \right) \left(-\frac{2}{M} k_2 - 2 \right) \right. \\ &\quad - t_{abc} I_{-1}^{a(1)} I_{-1}^{b(1)} I_{-1}^{c(2)} |0\rangle^{\mathbf{k}} \left(-\frac{2}{M} k_1 - 2 \right) \left(-\frac{2}{M} k_2 - 1 \right) \left(-\frac{2}{M} k_2 - 2 \right) \\ &\quad + t_{abc} I_{-1}^{a(1)} I_{-1}^{b(2)} I_{-1}^{c(2)} |0\rangle^{\mathbf{k}} \left(-\frac{2}{M} k_1 - 1 \right) \left(-\frac{2}{M} k_1 - 2 \right) \left(-\frac{2}{M} k_2 - 2 \right) \\ &\quad \left. - \frac{1}{3} t_{abc} I_{-1}^{a(2)} I_{-1}^{b(2)} I_{-1}^{c(2)} |0\rangle^{\mathbf{k}} \left(-\frac{2}{M} k_1 \right) \left(-\frac{2}{M} k_1 - 1 \right) \left(-\frac{2}{M} k_1 - 2 \right) \right), \end{aligned}$$

where the normalization factor $C(k_1, k_2)$ is given by

$$\begin{aligned} C(k_1, k_2) &:= -\frac{M^3}{4} \frac{1}{(k_1 + M)(k_2 + M)(k_1 + k_2 + M)} D(k_1, k_2), \\ D(k_1, k_2)^2 &:= \frac{-M}{2(M + 2k_1)(M + 2k_2)(3M + 2k_1 + 2k_2)(M^2 - 4)}. \end{aligned}$$

28 Sylvain Lacroix, Benoît Vicedo and Charles Young

Proposition 6.2. *Up to normalization factors depending only on the levels k_1, k_2 , the states ω and W are equal to $\int_{\gamma} \mathcal{P}(z)^{-1/M} \varsigma_1(z) dz$ and $\int_{\gamma} \mathcal{P}(z)^{-2/M} \varsigma_2(z) dz$ respectively. Namely,*

$$\omega = \frac{k_1 k_2}{(k_1 + k_2 + M)(k_1 + k_2)(k_1 + k_2 - M)} \frac{\int_{\gamma} \mathcal{P}(z)^{-1/M} \varsigma_1(z) dz}{\int_{\gamma} \mathcal{P}(z)^{-1/M} dz},$$

$$W = \frac{C(k_1, k_2)(-\frac{2}{M}k_1)_3(-\frac{2}{M}k_2)_3}{(-\frac{2}{M}k_1 - \frac{2}{M}k_2 + 1)_3} \frac{\int_{\gamma} \mathcal{P}(z)^{-2/M} \varsigma_2(z) dz}{\int_{\gamma} \mathcal{P}(z)^{-2/M} dz}.$$

Proof. From the definition, §3.3, of $\varsigma_1(z)$ we have, using Lemma 6.1,

$$(6.1) \quad \int_{\gamma} \mathcal{P}(z)^{-1/M} \varsigma_1(z) dz = - \int_{\gamma} \mathcal{P}(z)^{-1/M} \left(\frac{1}{z} - \frac{1}{z-1} \right) dz \times \Xi$$

where one finds

$$\Xi = -k_1 \omega^{(2)} - k_2 \omega^{(1)} + I_{-1}^{a(1)} I_{-1}^{a(2)} |0\rangle^{\mathbf{k}} = -(k_1 + k_2 + M) \omega.$$

The second equality is by a short calculation starting from the definition of ω above. By further use of Lemma 6.1, one has the first part. Similarly, from the definition of $\varsigma_2(z)$ in §3.3 we have

$$\begin{aligned} \int_{\gamma} \mathcal{P}(z)^{-2/M} \varsigma_2(z) dz &= \frac{1}{3} t_{abc} I_{-1}^{a(1)} I_{-1}^{b(1)} I_{-1}^{c(1)} |0\rangle^{\mathbf{k}} B(-\frac{2}{M}k_1 - 3, -\frac{2}{M}k_2) \\ &\quad + t_{abc} I_{-1}^{a(1)} I_{-1}^{b(1)} I_{-1}^{c(2)} |0\rangle^{\mathbf{k}} B(-\frac{2}{M}k_1 - 2, -\frac{2}{M}k_2 - 1) \\ &\quad + t_{abc} I_{-1}^{a(1)} I_{-1}^{b(2)} I_{-1}^{c(2)} |0\rangle^{\mathbf{k}} B(-\frac{2}{M}k_1 - 1, -\frac{2}{M}k_2 - 2) \\ &\quad + \frac{1}{3} t_{abc} I_{-1}^{a(2)} I_{-1}^{b(2)} I_{-1}^{c(2)} |0\rangle^{\mathbf{k}} B(-\frac{2}{M}k_1, -\frac{2}{M}k_2 - 3) \end{aligned}$$

and the second result follows by repeated use of Lemma 6.1. \square

By the coset construction [GKO86], the state $\omega \in \mathbb{V}_0^{\mathbf{k}}$ is a conformal vector: it generates a copy of the Virasoro vertex algebra with central charge

$$c = \dim(\mathfrak{sl}_M) \left(\frac{k_1}{k_1 + M} + \frac{k_2}{k_2 + M} - \frac{k_1 + k_2}{k_1 + k_2 + M} \right).$$

That is,

$$\omega_{(n)} \omega = \delta_{n,0} T(\omega) + \delta_{n,1} 2\omega + \delta_{n,3} \frac{1}{2} c |0\rangle^{\mathbf{k}}$$

for $n \in \mathbb{Z}_{\geq 0}$.

Cubic hypergeometric integrals of motion in affine Gaudin models 29

The state W was constructed in [BBSS88] (see equation 2.8 of that paper, where the field $Y[W, u]$ is given). Suppose we specialize further to the case of $M = 3$, *i.e.* $\widehat{\mathfrak{sl}}_3$, and choosing the irreducible module at the marked point $z_2 = 1$ to have highest weight

$$\lambda_2 = \Lambda_0.$$

so that, in particular, $k_2 = 1$. It was shown in [BBSS88] that in that case the states W and ω generate a copy of the W_3 algebra. That is, one has – see *e.g.* [BMP96] –

$$\omega_{(n)}W = \delta_{n,0}T(W) + \delta_{n,1}3W,$$

$$(6.2) \quad \begin{aligned} W_{(n)}W = \delta_{n,0} \left(\beta T(\Lambda) + \frac{1}{15}T^3(\omega) \right) + \delta_{n,1} \left(2\beta\Lambda + \frac{3}{10}T^2(\omega) \right) \\ + \delta_{n,2}T(\omega) + \delta_{n,3}2\omega + \delta_{n,5}\frac{1}{3}c|0\rangle^k, \end{aligned}$$

for $n \in \mathbb{Z}_{\geq 0}$. Here

$$\Lambda := \omega_{(-1)}\omega - \frac{3}{10}T^2(\omega), \quad \text{and} \quad \beta = \frac{16}{22 + 5c}.$$

Remark. The W_3 algebra is known to have a commutative algebra of Quantum Integral of Motions (of $\widehat{\mathfrak{sl}}_{3-(m)}$ -KdV, *i.e.* quantum Boussinesq theory) indexed by the exponents of $\widehat{\mathfrak{sl}}_3$. The first few integrals of motion I_1, I_2, I_4, I_5 can be found in [BHK02].³ The first two are $I_1 = \omega_0$ and $I_2 = W_0$. By Proposition 6.2, these are proportional to our first two Hamiltonians \hat{Q}_1^γ and \hat{Q}_2^γ . It is natural to think that the higher integrals of motion are likewise the two-point specializations of the (conjectural) higher Gaudin Hamiltonians. (This was suggested in the $\widehat{\mathfrak{sl}}_2$ case in [FF11, §6.4]). We have checked that the vacuum eigenvalues of I_4 and I_5 indeed agree with the eigenvalues $\int_\gamma \mathcal{P}(z)^{-4/3}v_4(z)dz$ and $\int_\gamma \mathcal{P}^{-5/3}(z)v_5(z)dz$ (see §4) predicted in [LVY], up to overall factors depending on the remaining weight λ_1 only through the level k_1 . The analogous check also works for I_3 and I_5 in the $\widehat{\mathfrak{sl}}_2$ case. \triangleleft

³Note that $W_{BBSS} = \sqrt{3\beta}W_{BHK}$.

References

- [BBSS88] F. A. Bais, P. Bouwknegt, M. Surridge, and K. Schoutens, *Coset construction for extended Virasoro algebras*, Nuclear Phys. B **304** (1988), no. 2, 371–391. MR 952975
- [BHK02] V. V. Bazhanov, A. N. Hibberd, and S. M. Khoroshkin, *Integrable structure of \mathcal{W}_3 conformal field theory, quantum Boussinesq theory and boundary affine Toda theory*, Nuclear Phys. B **622** (2002), no. 3, 475–547. MR 1880098
- [BLZ96] V. V. Bazhanov, S. L. Lukyanov, and A. B. Zamolodchikov, *Integrable structure of conformal field theory, quantum KdV theory and thermodynamic Bethe ansatz*, Comm. Math. Phys. **177** (1996), no. 2, 381–398. MR 1384140
- [BMP96] P. Bouwknegt, J. McCarthy, and K. Pilch, *The W_3 algebra*, Lecture Notes in Physics. New Series m: Monographs, vol. 42, Springer-Verlag, Berlin, 1996, Modules, semi-infinite cohomology and BV algebras. MR 1423803
- [FBZ04] E. Frenkel and D. Ben-Zvi, *Vertex algebras and algebraic curves*, second ed., Mathematical Surveys and Monographs, vol. 88, American Mathematical Society, Providence, RI, 2004. MR 2082709
- [FF11] B. Feigin and E. Frenkel, *Quantization of soliton systems and Langlands duality*, Exploring new structures and natural constructions in mathematical physics, Adv. Stud. Pure Math., vol. 61, Math. Soc. Japan, Tokyo, 2011, pp. 185–274. MR 2867148
- [FFR94] B. Feigin, E. Frenkel, and N. Reshetikhin, *Gaudin model, Bethe ansatz and critical level*, Comm. Math. Phys. **166** (1994), no. 1, 27–62. MR 1309540
- [Fre05] E. Frenkel, *Gaudin model and opers*, Infinite dimensional algebras and quantum integrable systems, Progr. Math., vol. 237, Birkhäuser, Basel, 2005, pp. 1–58. MR 2160841
- [Gau14] M. Gaudin, *The Bethe wavefunction*, Cambridge University Press, New York, 2014, Translated from the 1983 French original by Jean-Sébastien Caux. MR 3469644
- [GKO86] P. Goddard, A. Kent, and D. Olive, *Unitary representations of the Virasoro and super-Virasoro algebras*, Comm. Math. Phys.

Cubic hypergeometric integrals of motion in affine Gaudin models 31

103 (1986), no. 1, 105–119. MR 826859

- [Kac90] V. G. Kac, *Infinite-dimensional Lie algebras*, third ed., Cambridge University Press, Cambridge, 1990. MR 1104219
- [LVY] S. Lacroix, B. Vicedo, and C. A. S. Young, *Affine gaudin models and hypergeometric functions on affine opers*, preprint.
- [Mol13] A. I. Molev, *Feigin-Frenkel center in types B, C and D*, Invent. Math. **191** (2013), no. 1, 1–34. MR 3004777
- [MTV06] E. Mukhin, V. Tarasov, and A. Varchenko, *Bethe eigenvectors of higher transfer matrices*, J. Stat. Mech. Theory Exp. (2006), no. 8, P08002, 44. MR 2249767
- [Nah91] W. Nahm, *A proof of modular invariance*, Internat. J. Modern Phys. A **6** (1991), no. 16, 2837–2845, Topological methods in quantum field theory (Trieste, 1990). MR 1117751
- [Ryb] L. G. Rybnikov, *A proof of the Gaudin Bethe Ansatz conjecture*, preprint.
- [Ryb08] ———, *Uniqueness of higher Gaudin Hamiltonians*, Rep. Math. Phys. **61** (2008), no. 2, 247–252. MR 2424092
- [Tal06] D. V. Talalaev, *The quantum Gaudin system*, Funktsional. Anal. i Prilozhen. **40** (2006), no. 1, 86–91. MR 2223256
- [Vic] B. Vicedo, *On integrable field theories as dihedral affine Gaudin models*, preprint.

UNIV LYON, ENS DE LYON, UNIV CLAUDE BERNARD,
CNRS, LABORATOIRE DE PHYSIQUE, F-69342 LYON, FRANCE
E-mail address: sylvain.lacroix@ens-lyon.fr

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF YORK,
YORK YO10 5DD, U.K.
E-mail address: benoit.vicedo@gmail.com

SCHOOL OF PHYSICS, ASTRONOMY AND MATHEMATICS,
UNIVERSITY OF HERTFORDSHIRE, COLLEGE LANE, HATFIELD AL10 9AB, UK.

32 Sylvain Lacroix, Benoît Vicedo and Charles Young

E-mail address: `c.a.s.young@gmail.com`